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Article

Validating a Reduced-Order Model for Synthetic Jet Actuators Using CFD and Experimental Data

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Abstract: Synthetic jet actuators (SJA) are emerging in various engineering applications, from flow separation and noise control in aviation to thermal management of electronics. A SJA oscillates a flexible membrane inside a cavity connected to a nozzle producing vortices. A complex interaction between the cavity pressure field and the driving electronics can make it difficult to predict performance. A reduced-order model (ROM) has been developed to predict the performance of SJAs. This paper applies this model to a canonical configuration with applications in flow control and electronics cooling, consisting of a single SJA with a rectangular orifice, emanating perpendicular to the surface. The practical implementation of the ROM to estimate the relationship between cavity pressure and jet velocity, jet velocity and diaphragm deflection and applied driving voltage is explained in detail. Unsteady Reynolds-averaged Navier Stokes computational fluid dynamics (CFD) simulations are used to assess the reliability of the reduced-order model. The CFD model itself has been validated with experimental measurements. The effect of orifice aspect ratio on the ROM parameters has been discussed. Findings indicate that the ROM is capable of predicting the SJA performance for a wide range of operating conditions (in terms of frequency and amplitude).

Keywords: synthetic jet; reduced-order model; lumped-element model; gas dynamics; Helmholtz resonance; piezoelectric actuator; electromagnetic actuator; computational fluid dynamics

1. Introduction

Synthetic jet actuators (SJAs) generate a train of vortices which are formed by periodic suction and ejection of the same amount of fluid across an orifice. This establishes a directional flow with zero net mass input, formed from the ambient fluid surrounding the orifice, hence “synthetic” jet [1,2]. These types of actuators have been used to control flow separation and noise, e.g., in aviation applications, as well as more recently in the thermal management of electronics.

An SJA typically consists of a piezoelectric or electromagnetic driver deflecting a flexible membrane inside a cavity, which is connected to an external body of fluid through a short nozzle or orifice (The terms ‘nozzle’ and ‘orifice’ will be used interchangeably in this paper). An oscillation of the internal membrane causes a cavity pressure fluctuation and, thus, an oscillatory flow through the orifice. At a sufficiently high stroke length [3], a stable vortex detaches from the orifice and propagates into the surrounding flow field, thereby imparting momentum to a boundary layer.

A basic synthetic jet flow emanating into quiescent ambient fluid is characterized by a few parameters: (i) the jet Reynolds number $Re = U_0 d_h / \nu$, where d_h is the hydraulic diameter of the orifice, ν is the fluid kinematic viscosity and U_0 is a characteristic velocity scale representing the average

nozzle ejection velocity, defined as $U_0 = L_0 f$, where f is the actuation frequency. The stroke length L_0 is defined as the average distance a fluid slug is expelled during the ejection phase, assuming a harmonic velocity profile in the orifice which is positive during ejection ($0 \leq t < 0.5f^{-1}$) and negative during the suction phase ($0.5f^{-1} \leq t < f^{-1}$):

$$L_0 = \int_{t=0}^{0.5f^{-1}} u_n(t) dt \quad (1)$$

where $u_n(t)$ represents the spatially averaged orifice or nozzle velocity, hence subscript ‘n’. The dimensionless stroke length L_0/d_h is the reciprocal of a Strouhal number or $L_0/d_h = (fd_h/U_0)^{-1}$. Usually the actuator diaphragm is excited by a sine wave, in which case the characteristic velocity $U_0 = \pi^{-1}U_n$ where U_n is the peak (spatially averaged) orifice velocity, or the amplitude of the sine wave $u_n(t) = U_n \sin 2\pi ft$.

For sharp-edged orifices, threshold values for a vortex to detach, and thus for a synthetic jet flow to form, are $L_0/d_h > 0.50$ for a circular orifice and $L_0/d_h > 0.95$ for a two-dimensional slot orifice with infinite aspect ratio $a \rightarrow \infty$, where a is defined as the span-to-width ratio of a rectangular slot or $a = b/h$. The stroke length is a characteristic for the flow structure [4].

For flow control purposes, the SJA is usually embedded within a streamlined or bluff body, with the jet emanating from a surface at an angle [5,6]. For electronics cooling purposes, the synthetic jet flow is usually directed at a surface, impinging onto the object to be cooled [7]. In both cases however, the orifice can be circular or rectangular in shape, and the jet can issue at an oblique angle. Moreover, the use of dual adjacent SJAs can be used to achieve flow vectoring, determined by the phase difference in the driving signals for adjacent actuators [8].

Each of these configurations leads to quite different and complicated pressure fields, making it hard to predict SJA performance a priori. Since the jet flow structure is determined by the aforementioned dimensionless groups Re and L_0/d_h , information is needed about the synthetic jet orifice velocity $u_n(t)$, or for sine wave actuation, simply the velocity amplitude $U_n (= \pi U_0)$. Hence:

$$\begin{cases} Re = \frac{d_h}{\pi \nu} U_n \\ \frac{L_0}{d_h} = \frac{1}{\pi d_h} \frac{U_n}{f} \end{cases} \quad (2)$$

However, the orifice velocity cannot be inferred directly from the driving voltage e . Nor can the orifice velocity be measured easily in typical applications, requiring either intrusive techniques (e.g., hot-wire anemometry) or more expensive optical velocity measurements. A calibration correlating actuator voltage e with orifice velocity U_m can be performed; however, this may be subject to degradation or drift in actuator characteristics. A calibration of cavity acoustic pressure p_c with orifice velocity is, therefore, preferred [9–11]. For instance, when changing the phase difference between adjacent SJAs, only a pressure-velocity calibration approach can maintain a constant jet Reynolds number [8].

For a synthetic jet issuing into an external boundary layer, the combined flow field is further characterized by the velocity ratio U_p/U_∞ or the dimensionless frequency $F^+ = fL/U_\infty$, where U_∞ is the freestream velocity and L represents the streamwise location of the SJA from a leading edge, or the length of a separated zone.

Three main types of actuating mechanisms have been used to generate synthetic jet flows in academic research, (i) an oscillating piston, (ii) an electromagnetic loudspeaker or (iii) a piezoelectric bending element. Both air and liquid have been used as working fluids. It will be shown in Section 2 that the type of driver (i, ii or iii) does not affect the relationship between the nozzle velocity U_n and cavity pressure p_c , only the relationship between input power and diaphragm motion. Hence, a general reduced-order model (ROM) will incorporate a so-called ‘fluidic model’ relating U_n to p_c which does not depend on the actuator type, and a ‘driver model’ relating driver input to diaphragm deflection, x_d . The cavity pressure p_c is the physical linking variable between both parts of the model.

Table 1 gives an overview of different types of actuators, orifices, and SJA operating conditions for a selection of experimental studies in the scientific literature. The list is not exhaustive, since the purpose is merely to convey the range of actuator types and typical parameter ranges encountered. The final table entry represents the current study.

The objectives of this paper are threefold: (i) to provide a clear but concise overview of an analytical ROM for a generic SJA, which would be useful both for early stage design and for controlling SJA operation; (ii) to validate this ROM for a typical lab-scale SJA with rectangular orifice, using both experimental calibration data as well as 2D and 3D transient CFD simulation results; (iii) to compare CFD and experimental results and illustrate the potential challenges associated with the comparison of 2D and 3D geometries for SJA in the frequency domain.

Table 1. List of experimental synthetic jet investigations with details of the type of orifice and actuator used.

Study	Orifice	Fluid	Actuator Type	L_0/d_h	Re	U_0 , m/s	f , HZ
Smith and Glezer [10]	Slot ($h = 0.5$ mm, $a = 147$)	Air	Piezoelectric disk	14.6	596	8.5	577
Shuster and Smith [12]	Circular ($d_h = 25.4$ mm)	Water	Oscillating piston	1.0–3.0	1000–10,000	0.04–0.39	1.6–5.2
Smith and Swift [13]	Slot ($h = 5.1$ mm, $a = 47$)	Air	Loudspeakers	6.9–41	1360–28,790	2.0–41.8	29–102
Crittenden and Glezer [14]	Circular ($d_h = 1.6$ mm–4.8 mm)	Air	Oscillating piston	>76	989–35,830	9.0–108	25–200
Kordik and Trávníček [15]	Circular ($d_h = 10$ mm)	Air	Loudspeaker	6.3–15.8	2400–7100	3.5–10.3	55–65
Current study	Slot ($h = 1.5$ mm, $a = 30$)	Air	Loudspeaker	0.5–40	80–1836	0.2–9.5	41–164

2. Analytical Reduced-Order Model (ROM)

The SJA ROM used in this paper is a lumped parameter analytical model, combining (i) a second-order equivalent representation of the driving element (i.e., a piezoelectric diaphragm or a loudspeaker) with (ii) a zeroth-order gas dynamics approximation of the fluid motion in the orifice, with fluid compressibility in the cavity. The main mechanical degrees of freedom represent the motion of the diaphragm U_d and fluid slug in the nozzle U_n , respectively.

The overall system response function can be broken down into its constituent parts as follows:

$$\frac{\text{Fluidic power output}}{\text{Electrical power input}} \propto \frac{\rho c A_n U_n^2}{e^2 Z_e} = \underbrace{\frac{\rho c U_n}{p_c} \frac{A_n U_n}{A_d U_d} \frac{p_c A_d}{U_d}}_{\substack{\text{Operational} \\ \text{fluidic} \\ \text{model} \\ \text{Fluidic effects} \\ \text{in cavity and nozzle}}} \cdot \underbrace{\left(\frac{U_d}{e} \right)^2 Z_e}_{\text{Driver electro-mechanics}} \quad (3)$$

where ρ and c are the fluid density and speed of sound, A_n and A_d are the cross-section and surface areas of the nozzle and diaphragm, respectively, and Z_e is the electrical impedance of the driver. All variables in bold (U_n , U_d , e , p_c , Z_e) in Equation (3) represent complex quantities in the frequency domain, dependent on the Laplace variable $s = j2\pi f$ which has been omitted in Equation (3) for brevity.

2.1. Brief Description of the ROM

The operational fluidic model $\rho c U_n / p_c$ was described by Persoons and O'Donovan [11] and can be used to estimate the SJA velocity from a measured cavity pressure. It can be simplified to a closed form analytical expression, and thus is easily implemented in a data acquisition or control system, as described in this section. The coupling between the electromechanical driver model and fluidic model is described in full detail by Persoons [16]. Comparable reduced-order models, also referred to as lumped element models (LEM), are presented and discussed by other researchers [17–21]. Thus for

the sake of brevity and clarity, only the basic elements are described here and the reader is referred to other works for more details on the derivations of the equations [11,16].

Figure 1 shows a schematic diagram and equivalent electrical networks for this ROM with a piezoelectric actuator (Figure 1b) and electromagnetic (e.g., loudspeaker) actuator (Figure 1c). The diaphragm and nozzle slug velocities and displacements are related as $U_d = sx_d$ and $U_n = sx_n$.

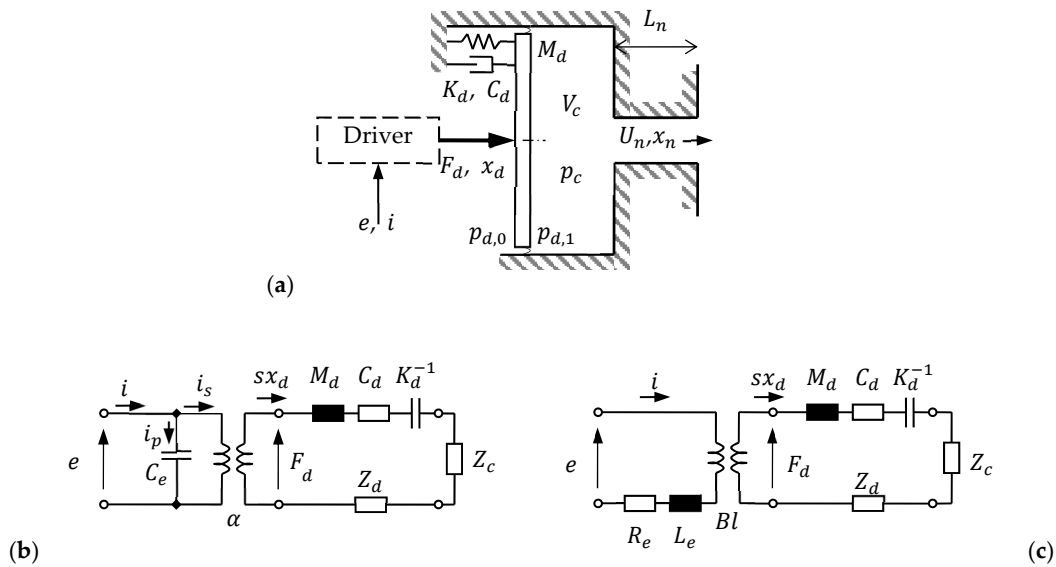


Figure 1. (a) Schematic diagram and equivalent electrical networks of a reduced-order model of a synthetic jet actuator (SJA) with (b) piezoelectric and (c) electromagnetic driver: from driving voltage e to diaphragm deflection x_d , cavity pressure p_c and nozzle velocity U_n [16].

Most parameters can be determined from known dimensions and manufacturer's specifications for the driver. Only two parameters (K and β) should be determined by calibration: K relates to the non-linear nozzle damping and β relates to the added nozzle inertia due to acoustic radiation.

The mechanical/fluidic behavior can be described by combining two second-order models, one relating to the oscillating diaphragm and one relating to the compressible gas dynamics:

$$\text{Diaphragm : } (M_d s^2 + C_d s + K_d) x_d = F_d - (p_{d,1} - p_{d,0}) A_d \quad (4)$$

$$\text{Gas dynamics : } (\rho A_n L_n s^2 + (C_n + Z_{n,1} + Z_{n,0}) s) x_n = p_c A_n \quad (5)$$

where M_d, C_d, K_d are the equivalent mass, damping and stiffness of the driver, F_d is the motive force applied by the actuator and $p_{d,1}$ and $p_{d,0}$ are the pressures on either side of the diaphragm, which are governed by acoustic radiation impedances Z_d and the cavity impedance Z_c . Equation (4) can be thus rewritten as:

$$\text{Diaphragm : } (M_d s^2 + (C_d + Z_c + Z_d) s + K_d) x_d = F_d = \alpha e \quad (6)$$

where α is an electromechanical coupling coefficient. In case of a piezoelectric actuator, $\alpha = d_{33} K_d$ where d_{33} is the piezoelectric modulus relating electric field and volumetric extension. This case is represented in Figure 1b.

In Equation (5), L_n and A_n are the geometric length and cross-sectional area of the nozzle, $C_n U_n$ represents the damping force related to fluid oscillating through the orifice, and the two Z_n terms correspond to acoustic impedances on both ends of the nozzle. C_n and Z_n are directly related to the aforementioned parameters K and β , as explained below.

Firstly, the frequency-dependent acoustic radiation impedance Z_n is described in detail by Beranek [22] for some canonical geometries. Appendix A in Persoons [16] describes the radiation

impedance for a circular and rectangular orifice of aspect ratio $a = 32$, based on tabulated data by Burnett and Soroka [23]. A low-frequency approximation for $f < c/(50A_n^{1/2})$ of this impedance $Z_n \cong M_f s = \rho A_n (\beta \sqrt{4A_n/\pi}) s$ where M_f represents a mass of fluid adjacent to one side of the orifice which adds to the oscillating inertia. This approximation thus leads to the introduction of the added mass coefficient β . When applied to both sides of the orifice equally, this can be recast into an equivalent (end-corrected) nozzle length L'_n :

$$L'_n = L_n + 2\beta \sqrt{4A_n/\pi} \quad (7)$$

With the total equivalent (quantities with a prime (') indicate the inclusion of an added mass correction) oscillating fluid mass $M'_n = \rho A_n L'_n$ and the cavity compressibility $K_c = \rho c^2 A_n^2 / V_c$ where V_c is the cavity volume, the Helmholtz resonance frequency emerges from the Equation:

$$f_H = \frac{1}{2\pi} \frac{c}{L'_n} \sqrt{\frac{A_n L'_n}{V_c}} \quad (8)$$

Secondly, the nozzle damping force $C_n \mathbf{U}_n$ can be represented by a combination of first-order (viscous) and second-order (inertial) damping. The relative dominance of first or second-order damping depends on the Stokes number ($\propto \sqrt{f d_h^2 / \nu}$); however, for short nozzles and moderate frequencies the first-order term can typically be ignored [16], leaving only:

$$C_n \mathbf{U}_n = K A_n \frac{\rho |\mathbf{U}_n|}{2} \mathbf{U}_n \quad (9)$$

Equation (5) can thus be recast as the non-linear cavity impedance Z_c relating SJA cavity pressure to orifice velocity:

$$Z_c = \frac{p_c}{\rho c \mathbf{U}_n} = \frac{\rho A_n L'_n s + \frac{1}{2} K \rho A_n |\mathbf{U}_n|}{\rho c A_n} \quad (10)$$

The reciprocal of Z_c in Equation (10) describes the operational SJA model $\rho c \mathbf{U}_n / p_c$ in the frequency domain. Appendix C in Persoons [16] explains how this expression can be simplified for sine wave actuation to the following closed-form expression to estimate the nozzle velocity amplitude U_n from the cavity pressure amplitude p_c :

$$\frac{\rho c U_n}{p_c} = \sqrt{\frac{2V_c}{A_n L'_n}} \frac{1}{\sqrt{\left(\frac{\omega}{\omega_H}\right)^2 + \sqrt{\left(\frac{\omega}{\omega_H}\right)^4 + K^2 \left(\frac{V_c}{A_n L'_n}\right)^2 \left(\frac{p_c}{\rho c^2}\right)^2}}} \quad (11)$$

Using a microphone, the cavity pressure amplitude p_c can easily be measured [11,16] and Equation (11) thus allows for U_n to be directly calculated, once values for K and β are known for the orifice. The following section describes the procedure for determining these parameters.

Other useful expressions relating nozzle velocity to diaphragm deflection x_d and actuator voltage e can be derived from the above set of equations. The reader is referred to Persoons [16] for the full derivation. The following relates nozzle velocity \mathbf{U}_n to diaphragm deflection $\mathbf{x}_d (= \mathbf{U}_d s^{-1})$:

$$\frac{\mathbf{U}_n A_n}{\mathbf{U}_d A_d} \left(= \frac{\mathbf{U}_n A_n}{s \mathbf{x}_d A_d} \right) = \frac{K_c}{M'_n s^2 + C_n s + K_c} \quad (12)$$

Based on derivations in Persoons [16], the relationship between nozzle velocity \mathbf{U}_n and piezoelectric actuator voltage e is given by:

$$\frac{\mathbf{U}_n}{e} = \frac{A_d}{A_n} \underbrace{\left(\frac{K_c}{M'_n s^2 + C_n s + K_c} \right)}_{\text{Fluidic effects in cavity and nozzle}} \underbrace{\left(\frac{\alpha s}{M_d s^2 + (C_d + Z_c + Z_d)s + K_d} \right)}_{\text{Piezoelectric actuator electromechanics}} \quad (13)$$

For an electromagnetic actuator (see Figure 1c), a similar expression can be obtained [16]:

$$\frac{\mathbf{U}_n}{e} = \frac{A_d}{A_n} \underbrace{\left(\frac{K_c}{M'_n s^2 + C_n s + K_c} \right)}_{\text{Fluidic effects in cavity and nozzle}} \underbrace{\left(\frac{(Bl)s / (R_e + L_e s)}{M_d s^2 + (C_d + Z_c + Z_d)s + K_d + K_{emf}} \right)}_{\text{Loudspeaker actuator electromechanics}} \quad (14)$$

where R_e and L_e are the voice coil resistance and inductance, respectively, Bl is the electromagnetic force factor ($F_d = Bl \cdot i$), and $K_{emf} = (Bl)^2 s / (R_e + L_e s)$.

Equations (12)–(14) are only provided here in complex notation, yet these can easily be implemented in a data acquisition software such as NI LabVIEW or MATLAB (sample code is available upon request from the corresponding author (Tim Persoons)) to infer the amplitude and phase lag of \mathbf{U}_n as a function of either x_d or e .

The two system resonance frequencies f_1 and f_2 are related to, but not equal to, the Helmholtz resonance frequency f_H and driver resonance frequency $f_d = \frac{1}{2\pi} \sqrt{K_d / M'_d}$, and are determined as:

$$\begin{aligned} f_1 &\cong \frac{1}{2\pi} \sqrt{\frac{[K_c^{-1} + ((A_n / A_d)^2 K_d)^{-1}]^{-1}}{M'_n}} \\ f_2 &\cong \frac{1}{2\pi} \sqrt{\frac{K_d + (A_d / A_n)^2 K_c}{M'_d}} \end{aligned} \quad (15)$$

The equivalent diaphragm mass appears here in its modified form [16], including an added mass correction on one side of the diaphragm, $M'_d = M_d + \beta \rho A_d (4A_d / \pi)^{1/2}$.

Persoons [16] gives more details about the physical interpretation of f_1 and f_2 , and further confirmation for these resonance frequencies can also be found in Gallas et al. [18], Kordik et al. [19] and Kooijman and Ouweltjes [24].

2.2. Estimating ROM Parameters K and β

The non-linear damping coefficient K and added mass coefficient β can be determined by least-squares fitting the model expression in a calibration procedure, as described by Persoons [16]. This will tend to over predict β as the fit includes acoustic effects from the diaphragm, which are only implicitly included in the ROM while only the acoustic effects at the nozzle are explicitly used. Alternatively, β can be determined analytically, as described below.

As mentioned in the brief description of the SJA ROM, β arises from acoustic radiation impedance and affects the inertia of the nozzle in the same manner that adding additional mass would. For a rectangular nozzle, such as the one used in this work, acoustic impedance can be interpolated from the tabulated data given by Burnett and Soroka [23] or analytically determined from the equation provided therein [23] and reproduced below. For a circular nozzle, the acoustic impedance can be determined by using a Struve H function of the first kind [16]. For a rectangular nozzle, β can be determined as a function of frequency and aspect ratio a as follows:

$$\beta = \frac{\sqrt{\pi} \chi(k A_n^{1/2}, a)}{2k A_n^{1/2}} \quad (16)$$

where χ is the imaginary part of the acoustic impedance and k is the wave number ($k = 2\pi f/c$) which is normalized by the nozzle cross-sectional area as $kA_n^{1/2}$. The function $\chi(\gamma, a)$ can be determined as:

$$\chi(\gamma, a) = \frac{2}{\pi\gamma^2} \left[\sin \gamma q - \gamma q \cos \gamma q + \gamma \left(p + \frac{1}{p} \right) - \sin \gamma p - \sin \frac{\gamma}{p} \right] - \frac{2}{\pi} \left[pI(\gamma, a) + \frac{I(\gamma, 1/a)}{p} \right] \quad (17)$$

where parameters p and q are defined as $p = a^{1/2}$ and $q = (a + 1/a)^{1/2}$, and the function $I(\gamma, \xi)$ is defined as:

$$I(\gamma, \xi) = \int_{\xi^{-1/2}}^{(\xi+1/\xi)^{1/2}} \left(1 - \frac{1}{\xi t^2} \right)^{1/2} \sin \gamma t dt \quad (18)$$

Note that Equations (16)–(18) for determining β are a low-frequency approximation, and only valid for values of $kA_n^{1/2} < 0.886$. Alternatively, the following equation is an empirical fit for β that is valid for $1 < a < 50$ and $kA_n^{1/2} < 0.886$ with a mean error of 1.7% and a maximum error of 5.9%:

$$\beta \approx -0.34064 + 6.6771 \times 10^{-2} kA_n^{1/2} + 1.5318 \times 10^{-4} a^2 + 2.7233 \times 10^{-8} a^4 + 0.78343 \exp(-1.7384 \times 10^{-2} a) \quad (19)$$

While Equation (19) is only valid for aspect ratios not exceeding 50, the value of β from Equations (16)–(18) in the limit for infinite aspect ratio ($a \rightarrow \infty$) is zero. Thus, for a two-dimensional CFD model with an inherent aspect ratio of infinity, β tends to zero and thus no end correction is needed for a ROM representing that case. In Equation (7), the apparent nozzle length L'_n for a 2D CFD simulation thus reduces to the geometric length L_n , which is used in Section 3.3 to explain how to compare the performance of 2D and 3D SJA cavity geometries.

3. Validation Methodologies

3.1. Numerical Validation Using Transient Computational Fluid Dynamics (CFD) Modelling

Two-dimensional (2D) unsteady compressible CFD simulations are carried out using Ansys CFX software. For a selected number of cases, three-dimensional (3D) simulations have been carried out using the same package (see Section 3.1.1). Figure 2 displays the 2D computational domain and the generated mesh used in the simulations. To achieve the most realistic computation of the flow induced by the synthetic jet, Alimohammadi et al. [25] reported on the importance of diaphragm deformation and the inclusion of the cavity and deforming diaphragm as a part of the simulation. This approach stands in contrast to the traditional assumption of an oscillating boundary condition at the exit of the orifice slot which leads to an over-simplification of boundary conditions. As a result, the oscillating wall of the cavity is simulated in the developed CFD model here by means of dynamic mesh techniques. The diaphragm deformation applied in the CFD model $x_d = f(t, y)$ is different from the deformation of the loudspeaker cone used for the experiments (see Section 3.2), however the maximum deflection for all cases considered in this study does not exceed 4% of the cavity height L_c , and therefore the difference in impact on the internal cavity flow is deemed negligible.

The computational domain consists of three distinct regions, namely cavity, nozzle (length $L_n = 10$ mm, width $h = 1.5$ mm), and the near-orifice region where the synthetic jet develops. The cavity dimensions for the simulations are $h_c = 75$ mm and $L_c = 23$ mm, to mimic the experimental validation geometry as best as possible.

The far-field boundaries at atmospheric pressure are placed at a distance of $120 h$ in the x direction and $\pm 70 h$ in the y direction from the orifice outlet. The simulation of synthetic jets consists of intricate fluid mechanics featuring flow separation, unsteadiness, and vortex dynamics. Alimohammadi et al. [25] have suggested the shear stress transport (SST) model coupled with modified curvature correction and also the Gamma-Theta transition model as the most suitable Reynolds Averaged Navier Stokes (RANS) turbulence model to be utilized. As such, the 2D domain is composed of nearly 100,000 cells in a structured mesh, and the near wall refinement assures the y^+ value (i.e., the non-dimensional distance of the first grid point from the wall) does exceed unity.

The numerical uncertainty is quantified by the grid convergence index (GCI) method which represents the discretization error, the method recommended by the *Journal of Fluids Engineering* [26]. Through several mesh refinements to study the independency of the mesh generation, the maximum numerical uncertainties (GCI_{max}) are calculated for different grids, reporting that solutions are within the asymptotic range of convergence. The GCI_{max} based on the peak nozzle velocity is calculated as 0.8%.

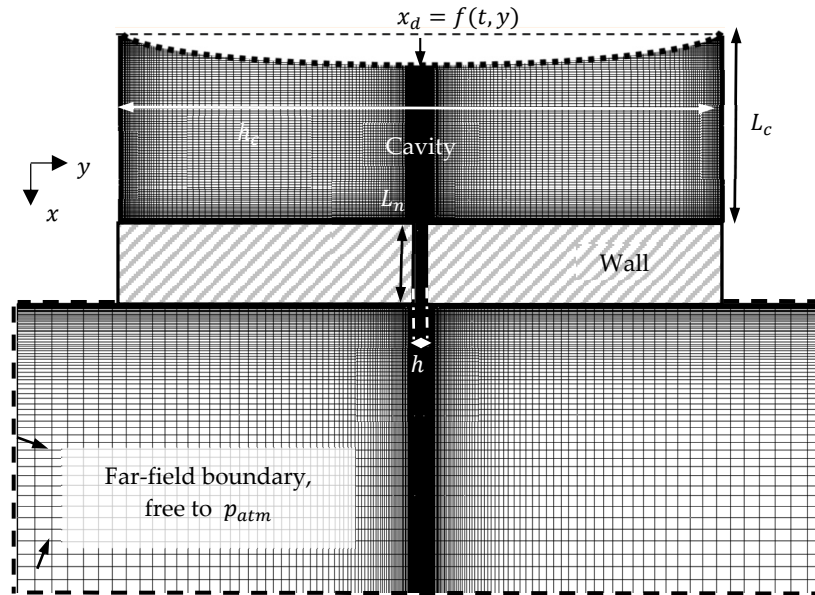


Figure 2. 2D computational domain and mesh (the diaphragm deformation is exaggerated and only indicative. The full solution domain is not shown for the sake of clarity).

3.1.1. Complementary 3D Simulations

For some selected cases, 3D CFD simulations are performed on a mesh with 1.46 million cells for the same solver settings and conditions. For the 3D simulations, the exact experimental dimensions are taken (see Section 3.2). By choosing similar linear cavity dimensions, it is anticipated that internal vortex formation in the cavity would be reasonably similar in 2D and 3D cases. However, Section 3.3 explains important consequences for the comparison of results from 2D and 3D geometries in the frequency domain.

3.2. Experimental Validation

The experimental validation approach is identical to the description in Persoons [16], to which the reader is referred for full details. Similar to the CFD simulations, the orifice slot measures $h = 1.5$ mm by $b = 45$ mm (aspect ratio $a = b/h = 30$) with a length $L_n = 10$ mm, while the cavity volume $V_c = 113$ cm³ and $A_d = 44$ cm², corresponding to a 75 mm circular diaphragm representing a Visaton FR8 (4 Ohm, 10 W) loudspeaker [16].

Figure 3 shows a schematic diagram identifying the primary instrumentation used: (i) a hot-wire sensor (e.g., Dantec 55P11, Pt-plated tungsten wire, 5 μ m diameter by 1.25 mm in length), operated in a constant temperature anemometer (Dantec 54T30, bridge ratio 20:1, resistance overheat ratio 1.8, 10 kHz low pass filter) measures the approximate centerline nozzle velocity $U_{n,0}(t)$ which is taken to represent the mean nozzle velocity in the assumption that the profile is reasonably uniform for a short orifice. The probe is less than 0.5 mm away from the orifice outlet plane to avoid velocity decay. The off-angle positioning of the probe, away from the main jet formation region, minimizes its flow disturbance. The hot-wire was calibrated in this same orientation in a low-turbulence wind tunnel, using a pitot-static probe as reference, for a range of air velocities between 0.6 m/s and 48 m/s.

A King's law relation was least-squares fitted to the calibration data, resulting in a coefficient of determination of 99.9%. The probe reading is insensitive to flow direction, yet due to the probe body orientation being nearly perpendicular to the jet flow (see Figure 3), the reading can be inverted during the suction stroke. The transitions between suction and ejection strokes are determined by locating the local minima in the hot-wire reading (in absolute value). During these transitions, wherever the velocity reading falls below the lower limit of the calibration range (<0.6 m/s), readings are ignored. (ii) The cavity acoustic pressure p_c is measured relative to atmosphere using a low sensitivity microphone (G.R.A.S., Hole, Demark. 40BD with 26CB constant current power (CCP) preamplifier, 1.6 mV/Pa, 40–174 dB, 4 Hz–70 kHz). (iii) The deflection of the diaphragm x_d is measured with a laser displacement sensor (e.g., Keyence LK-G157, range ± 40 mm, linearity $\pm 0.05\%$). These quantities are read into a data acquisition system along with the driver voltage e and current i . The data is phase-averaged over 16 periods to achieve an uncertainty level below 5% based on a 95% confidence level on the phase-averaged velocity, pressure and diaphragm deflection waveforms.

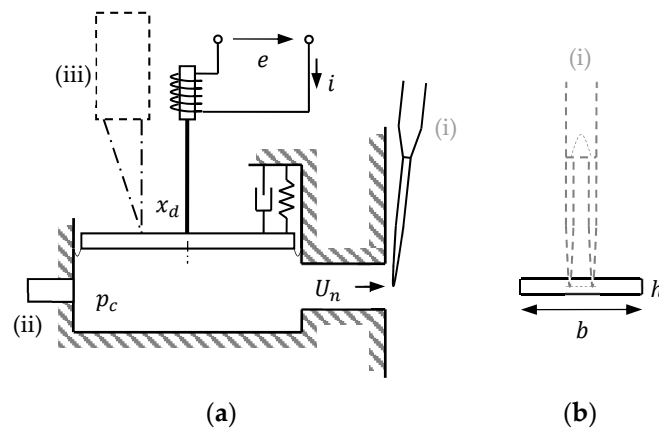


Figure 3. (a) Schematic diagram of an instrumented SJA with (b) rectangular orifice (*not to scale*), instrumented with (i) a hot-wire anemometer probe for nozzle velocity U_n , (ii) microphone for cavity pressure p_c and (iii) laser displacement sensor for x_d [16].

3.3. Comparing Synthetic Jet Actuator Performance for 2D and 3D Geometries

As mentioned in Section 2.1, the Helmholtz resonance frequency and driver resonance frequency are two important quantities that characterize the performance of an SJA in the temporal or frequency domain. The Helmholtz resonance frequency f_H defined by Equation (8) is the primary parameter, since typically the driver and cavity resonance frequencies are chosen to be of similar magnitudes for optimal performance.

For 3D geometries such as the actuator used for the experimental validation, the Helmholtz resonance frequency $f_H^{(3D)}$ can be estimated based on the cavity volume, nozzle cross-sectional area and end-corrected nozzle length (which requires an estimate of the added mass coefficient β):

$$f_H^{(3D)} = \frac{1}{2\pi} \frac{c}{L'_n} \sqrt{\frac{A_n L'_n}{V_c}} \text{ where } L'_n = L_n + 2\beta \sqrt{4A_n/\pi} \quad (20)$$

For a 2D geometry used for the CFD simulations in this study, the added mass coefficient $\beta = 0$, as explained in Section 2.2. Thus the 2D equivalent of the Helmholtz resonance frequency $f_H^{(2D)}$ is defined slightly differently as:

$$f_H^{(2D)} = \frac{1}{2\pi} \frac{c}{L'_n{}^{(2D)}} \sqrt{\frac{h L'_n{}^{(2D)}}{h_c L_c}} \text{ where } L'_n{}^{(2D)} = L_n \quad (21)$$

In this paper, the Helmholtz resonance frequencies for the 2D and 3D CFD simulations and experiments are $f_H^{(2D)} = 509$ Hz, $f_H^{(3D)} = 332$ Hz and $f_H^{(exp)} = 307$ Hz, respectively. In the following section, the frequency domain results are plotted as a function of the normalized frequency $\omega/\omega_H (= f/f_H)$, where f_H takes the aforementioned values for 2D and 3D data.

4. Results and Discussion

Table 2 shows an overview of the cases studied in this paper. All cases represent experiments and matching CFD simulations for an SJA with a rectangular orifice, however as specified above, there are inevitable physical differences in the systems, more specifically the 2D CFD geometry having an infinite orifice aspect ratio a , whereas the aspect ratio for the experiments and 3D CFD results is finite ($a = 30$).

Table 2 also summarizes the velocity-to-pressure ratio values, representing the fluidic SJA model defined in Equation (11). The values are provided in dimensionless form as $(\rho c U_n / p_c) (A_n L_n' / V_c)^{1/2}$, where the amplitudes U_n and p_c are obtained from the (i) experiments, (ii) 2D CFD and (iii) 3D CFD simulations, respectively. For each condition, the corresponding reduced-order model value obtained using Equation (11) is also provided in brackets.

Figure 4 provides a graphical representation of the agreement between the experimental values (markers) and the ROM (solid lines) in terms of the ratio of nozzle velocity to cavity pressure $\rho c U_n / p_c$, referred to as the fluidic model and defined in Equation (11). Because the model is non-linear, four lines are shown for four different cavity pressure amplitudes ranging from 100 Pa to 1000 Pa. The markers represent experimental measurement points at corresponding pressure amplitudes, as listed in Table 2. Cases A and B (see Table 2) are indicated with red and blue markers, respectively.

Table 2. List of cases included in this investigation. Cases A and B are indicated respectively as red and blue markers in subsequent figures. Values in parenthesis are reduced-order model (ROM) predictions at the operating conditions.

Case	Frequency $\frac{f}{f_H} (= \frac{\omega}{\omega_H})$	Pressure Amplitude p_c , Pa	Velocity-to -Pressure Ratio $(\frac{\rho c U_n}{p_c}) \sqrt{\frac{A_n L_n'}{V_c}}$, in dB		
			Experimental ($f_H^{(exp)} = 307$ HZ)	2D Computational Fluid Dynamics (CFD) ($f_H^{(2D)} = 509$ HZ)	3D CFD ($f_H^{(3D)} = 332$ HZ)
A	0.13	106	13.1 (13.0)	15.0 (15.1)	11.8 (13.0)
B	0.53	202	4.8 (5.1)	6.5 (9.1)	6.1 (5.5)
C	0.13	208	10.9 (10.7)	12.0 (12.3)	-
D	0.27	509	9.6 (9.1)	13.5 (12.5)	-
E	0.27	1027	6.8 (6.3)	8.7 (8.5)	-
F	0.27	202	3.5 (3.7)	6.5 (6.2)	-
G	0.53	504	4.6 (4.2)	5.5 (6.8)	-

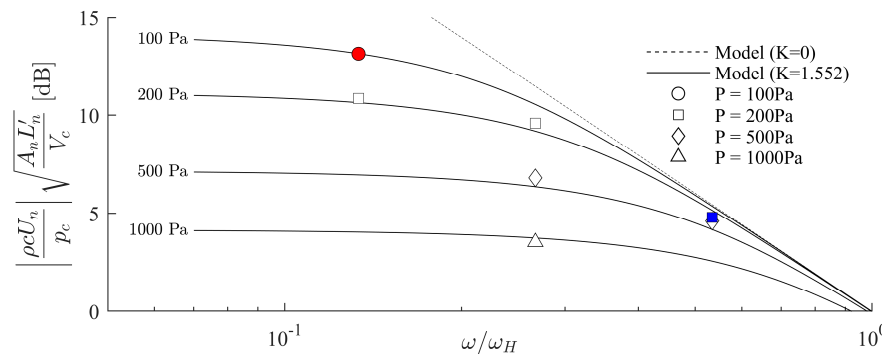


Figure 4. ROM validation in terms of the fluidic model U_n/p_c (Equation (11) with $K = 1.552$ and $\beta = 0.615$) (lines) versus experimental data (markers) for a synthetic jet actuator with rectangular orifice of aspect ratio 30:1. Markers represent cavity pressure magnitudes (\circ) 100 Pa, (\square) 200 Pa, (\diamond) 500 Pa, (\triangle) 1000 Pa. Cases A and B (see Table 2) are indicated with red and blue markers, respectively.

The higher the pressure amplitude, the more the damping term featuring K in Equation (11) dominates, and the less sensitive the velocity-to-pressure ratio becomes to frequency. The frequency on the horizontal axis is normalized by the Helmholtz resonance frequency. For frequencies exceeding the Helmholtz resonance by a factor of 3–5, the velocity-to-pressure ratio is no longer quasi constant and increasingly tends towards the dashed line. This represents an undamped model ($K = 0$), corresponding to $\frac{\rho c U_n}{p_c} = \sqrt{\frac{V_c}{A_n L_n}} \left(\frac{\omega}{\omega_H} \right)^{-1}$.

Regarding the procedure for determining the parameters K and β through calibration, this is described in detail in Persoons [16]. In short, once a given dataset of cavity pressure and nozzle velocity waveforms are provided, whether through experimental measurements or numerical simulations, the amplitudes U_n and p_c are determined to calculate the velocity-to-pressure ratio as listed in Table 2. Then a least-squares fitting procedure is applied to determine K and β such that Equation (11) best fits the experimental or numerically determined ratios. Because of the non-linearity, it is convenient for plotting purposes if the data are taken at predefined cavity pressure amplitudes, which explains why the data in Table 2 have cavity pressure amplitudes that are approximately 100 Pa, 200 Pa, 500 Pa and 1000 Pa. During the experiments, this can be done by monitoring the cavity pressure microphone readings, although admittedly this is somewhat more cumbersome to achieve for the CFD simulations using the dynamic meshing approach. However, this is merely a small inconvenience from the perspective of plotting results to collapse onto a small number of lines representing Equation (11).

The nozzle velocity, cavity pressure and diaphragm deflection waveforms that were used to determine the amplitudes in Figure 4 and Table 2 are shown for two selected cases A and B in Figures 5 and 6. These cases are chosen for the following reasons: Case A (red marker in Figure 4) is relatively low frequency and has a small cavity pressure amplitude (approximately 100 Pa), and thus also a small velocity amplitude, which makes for challenging measurement conditions. Case B (blue marker in Figure 4) is obtained at a higher frequency closer to the Helmholtz resonance frequency, which makes for increasingly more challenging conditions to obtain reliable convergence of the CFD simulations. Due to the significant computational cost involved in performing transient 3D CFD simulations, only cases A and B were simulated in 3D. 2D CFD simulations were carried out for the entire set of cases.

For Case A, Figure 5 shows plots of the phase-resolved (a) diaphragm deflection $x_d(\omega t)$, (b) cavity pressure $p_c(\omega t)$ and (c) nozzle velocity $U_n(\omega t)$ determined from experimental measurements, 2D CFD simulations and 3D CFD simulations, where the phase angle is given by $\omega t = 2\pi f t$. Phase angle $\omega t = 0$ is arbitrarily chosen as the zero crossing of the diaphragm deflection during the ejection stroke. For the experiments, the waveforms are ensemble-averaged over 16 cycles, as explained in Section 3.2. For the CFD simulations, the waveforms are extracted after periodic conditions were reached. The threshold for reaching periodicity is defined such that the root-mean-squared deviation on the instantaneous values of velocity and pressure between two successive cycles does not exceed 1%.

Since the diaphragm deflection is a user-defined input in the CFD simulations, perfect sine waveforms are observed for the CFD implementation of $x_d(\omega t)$ in Figure 5a. The simulated cavity pressure in Figure 5b broadly resembles the measured pressure waveform; however, some phase lag can be seen in some of the higher order fluctuations. With respect to the ROM, however, it only considers first order effects and neglects any higher order acoustics, thus these discrepancies are of minor importance in that regard. The difference in pressure magnitude in Figure 5b between experiments and CFD simulations has been verified to be independent of mesh density, turbulence model, and numerical discretization method. Furthermore, 2D and 3D CFD simulations yield similar pressure amplitudes. However, it is the ratio of velocity to pressure amplitude that matters for the fluidic model validation, not the cavity pressure amplitude itself.

In terms of the nozzle velocity $U_n(\omega t)$, it is important to note that the measured velocity is obtained by means of a hot-wire anemometer probe, placed in close proximity to the outlet of the jet nozzle (see Section 3.2). Because a hot-wire anemometer is insensitive to flow direction, the original velocity waveform is unfolded during the suction stroke to produce the directional velocity waveform

shown in Figure 5a. Furthermore, because of the asymmetry in the near-orifice flow fields for a synthetic jet during ejection and suction, there is a small but significant difference in the positive and negative peak velocity magnitudes during ejection and suction, respectively.

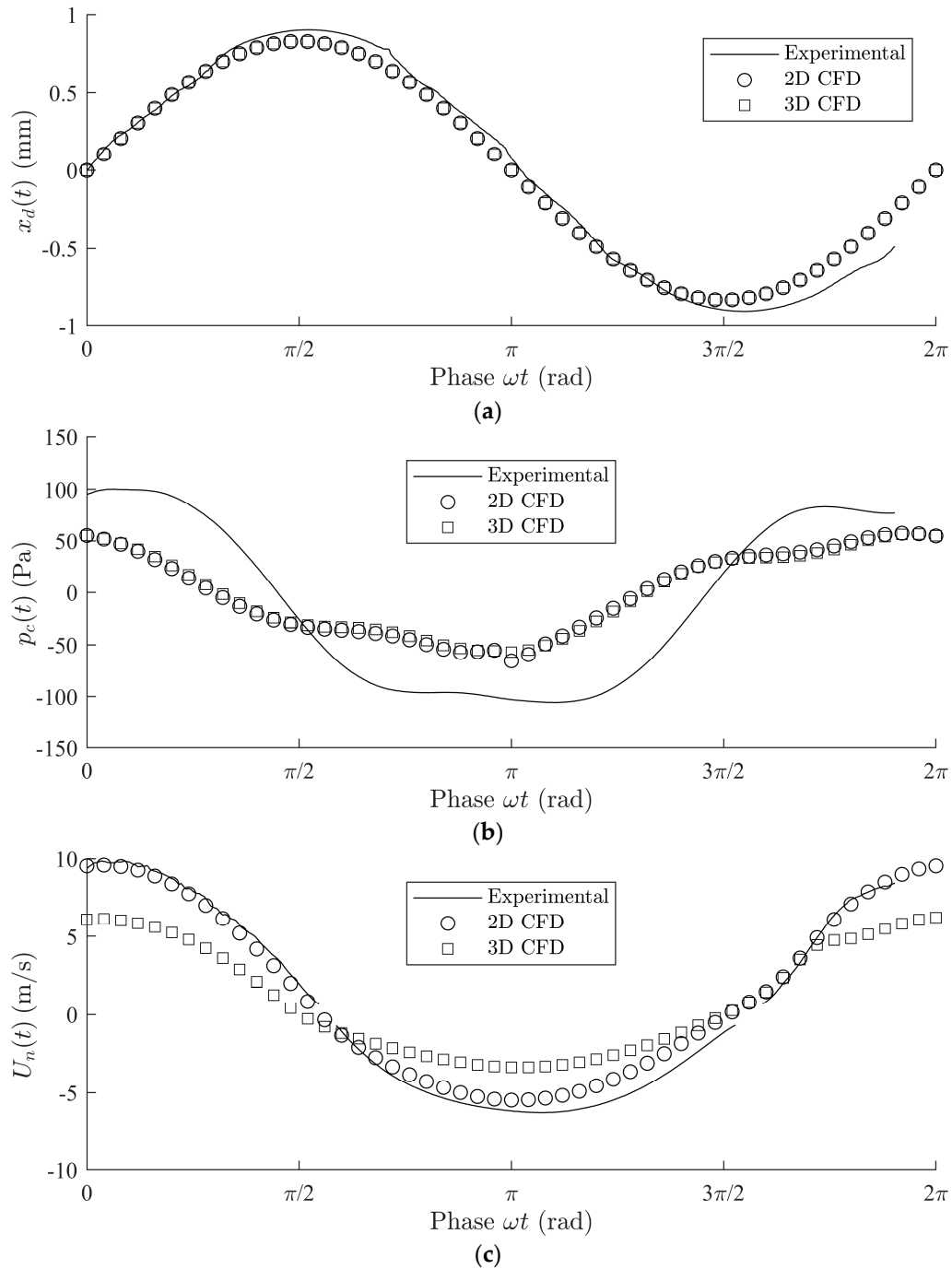


Figure 5. (a) Diaphragm deflection $x_d(\omega t)$, (b) cavity pressure $p_c(\omega t)$ and (c) nozzle velocity $U_n(\omega t)$ as a function of phase angle ωt for case A (see Table 2), comparing experimental results (—), 2D CFD results (○), and 3D CFD results (□).

Furthermore, as mentioned in Section 3.2, the hot-wire measures the approximate centerline nozzle velocity $U_{n,0}(\omega t)$. Therefore, the velocity at the same location is extracted from the CFD results to provide a fair comparison. In Figure 5c and 6c, hot-wire readings below the calibration range (<0.6 m/s) are omitted from the graphs.

Overall, Figure 5c shows a reasonable agreement between the experimental and numerical velocity waveforms, faithfully reproducing a minor peak during the startup of the ejection phase, visible at phase angle $\omega t / (2\pi) \cong 0.90$. The difference in magnitude between the experimental and CFD results should be interpreted as follows: although the 2D CFD results show a good agreement with the experiments in terms of velocity, the 2D CFD cavity pressure is underestimated. For the 3D CFD results, the predicted velocity magnitude is smaller yet the ratio of velocity to pressure amplitude is much closer to the experimental data, giving a stronger confirmation of the validity of the fluidic model in Equation (11).

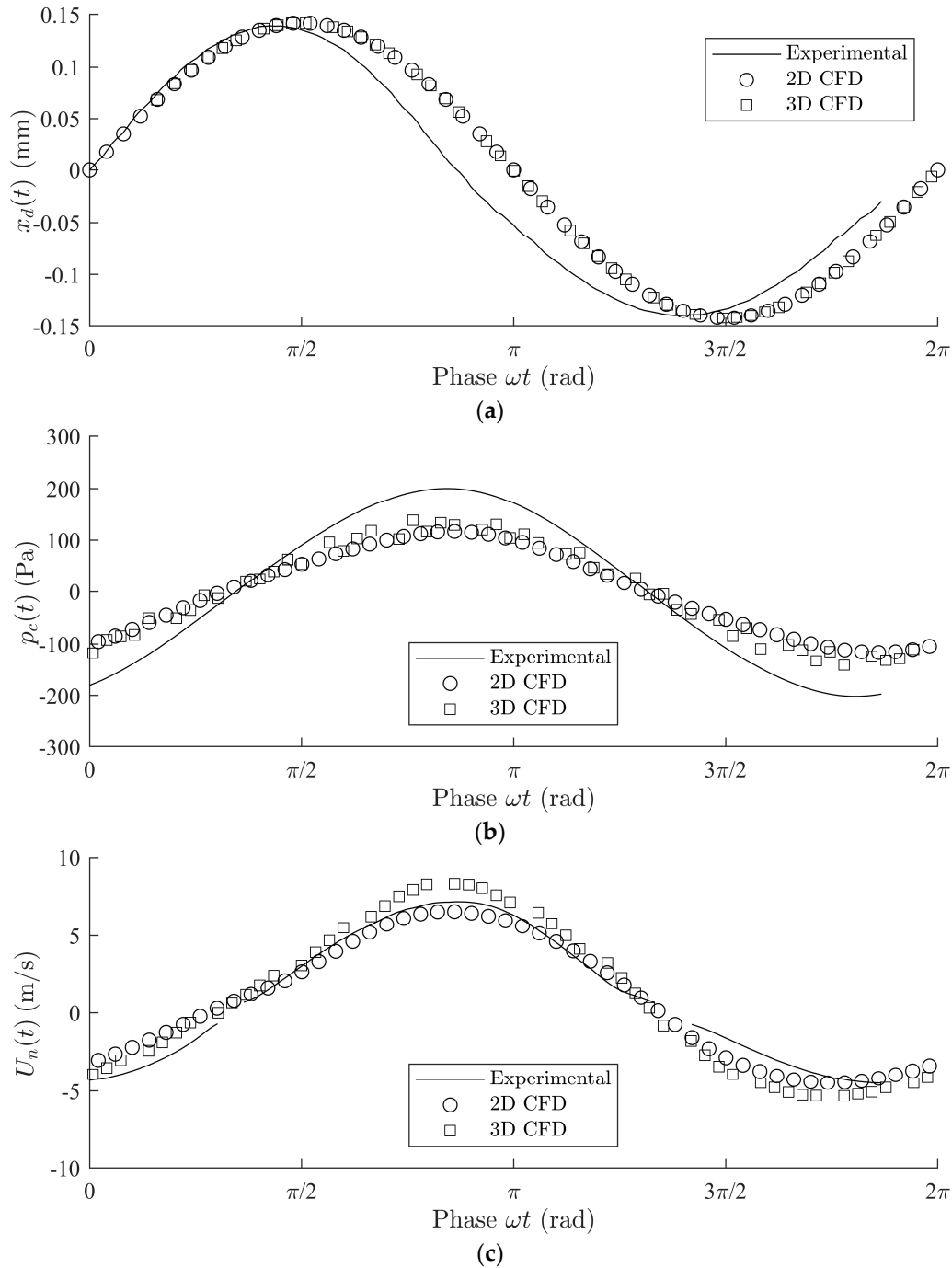


Figure 6. (a) Diaphragm deflection $x_d(\omega t)$, (b) cavity pressure $p_c(\omega t)$ and (c) nozzle velocity $U_n(\omega t)$ as a function of phase angle ωt for case B (see Table 2), comparing experimental results (—), 2D CFD results (○), and 3D CFD results (□).

Similar to Figure 5, Figure 6 shows the corresponding waveforms for case B at a greater frequency, closer to the Helmholtz resonance frequency. Again, a reasonable general agreement is observed between experimental and numerical data. Because it is approaching the resonance frequency, a significant phase shift is observed for case B between the diaphragm deflection and nozzle velocity, compared to the low frequency case A.

In terms of the 2D CFD simulations, it was found that the results are not particularly sensitive to changes in the modelling approach. Different discretization schemes and turbulence models were used for the 2D CFD simulations during the initial exploratory simulations for this study, and even the 3D CFD simulation results do not appreciably alter the shape of the pressure and velocity waveforms. A notable improvement in the velocity-to-pressure ratio prediction was observed for 3D CFD simulations compared to 2D simulations (see Figure 5).

Furthermore, it was found that there was no need to match exactly the location of the cavity pressure probe in the CFD pressure field to the location of the microphone used in the experiments. Thus it seems that, at least in first approximation, the physics relating bulk cavity pressure to nozzle velocity can indeed be captured reliably by a reduced-order model.

Figure 7 demonstrates this more clearly by plotting the experimental and numerical values for the velocity-to-pressure ratio on the vertical axis against the corresponding ROM prediction on the horizontal axis. As before, cases A and B are indicated by red and blue markers, respectively, and the other cases listed in Table 2 are also included (only 2D CFD and experimental data). As mentioned previously, because of the difference in geometry for the 2D versus 3D geometries, different β values are used. In Figure 7, the ROM predictions corresponding to the experimental and numerical data are produced using Equation (11) with $K = 1.552$. For the experimental results, $\beta = 0.615$. For the 3D CFD results, $\beta = 0.237$ and for the 2D CFD results, $\beta = 0$ as explained in Section 2.2. The reason for the difference in β between the 3D CFD and experimental results can be attributed to the different acoustic radiation impedance boundary conditions; the experiments feature radiation from the rear of the diaphragm, which is not included in the CFD simulation. Furthermore, the far field radiation and reflective properties of the lab environment are difficult to reproduce faithfully in CFD at a reasonable computational cost.

This study has revealed the challenges in reliably matching up experimental and numerical conditions for this type of phenomenon. Further research including more detailed pressure and velocity field quantification will be required to fully understand these aspects.

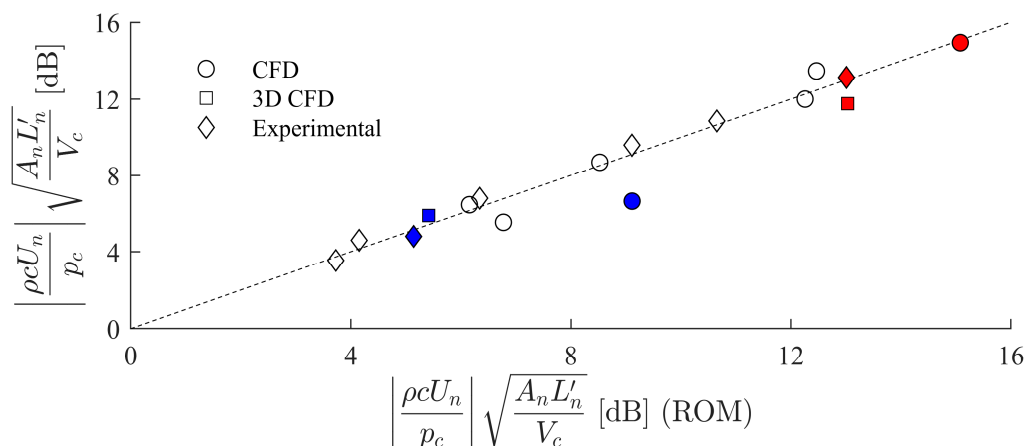


Figure 7. ROM validation in terms of the fluidic model U_n/p_c given by Equation (11) (horizontal axis) versus (○) 2D CFD results, (□) 3D CFD results and (◇) experimental results, for the same cases included in Figure 4 and Table 2. Cases A and B are indicated with red and blue markers, respectively.

5. Conclusions

In this paper, the underlying physics and methodology behind a physics-based reduced-order model (ROM) with two degrees of freedom for a synthetic jet actuator (SJA) has been explained. The ROM combines (i) a fluidic model relating jet velocity amplitude U_n to cavity pressure amplitude p_c and (ii) an electromechanical driver model relating the driver input to the diaphragm deflection amplitude x_d and cavity pressure amplitude p_c .

The fluidic model provides the velocity-to-pressure amplitude ratio U_n/p_c as a function of the fluid speed of sound, SJA dimensions and the operating frequency f . It is a non-linear model which accounts for fluid damping and inertia effects in the nozzle—these effects are captured by the two empirical model parameters, the pressure loss coefficient K and added mass coefficient β . Guidelines are provided for evaluating K and β , and a closed form expression for the added mass coefficient β is given for rectangular slot orifices as a function of orifice aspect ratio and frequency f .

The simplified version of the fluidic model equation (assuming sine wave actuation) is described by Equation (11) which is generally applicable to any SJA operating in any gas or liquid medium, as long as the fluid compressibility outweighs the mechanical compliance of the cavity and surrounding structure [27]. The electromechanical driver model is established in Equations (13) and (14) for the most common actuation methods used, i.e., a piezoelectric diaphragm and loudspeaker, respectively.

The paper has presented both experimental and numerical results using transient computational fluid dynamics (CFD) simulations, using both 2D and 3D geometries. Important aspects related to the comparison of 2D and 3D geometries are discussed in Sections 3.3 and 4.

A satisfactory agreement is achieved with the ROM predictions for both experimental measurements as well as 2D and 3D numerical simulations, as summarized by Figure 7. This demonstrates the validity of this fairly simple analytical ROM for a variety of conditions, for a range of frequencies approaching the Helmholtz resonance frequency.

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Nomenclature

A	Cross-sectional area, m^2
a	Orifice aspect ratio ($a = b/h$)
Bl	Electromagnetic force factor ($F_d = Bl \cdot i$)
b	Spanwise length of slot orifice, m
C	Linear damping coefficient, $\text{N}/(\text{m}^*\text{s})$ or capacitance, F
c	Speed of sound, m/s
d_h	Hydraulic diameter, m
e	Actuator voltage, V
F^+	Dimensionless frequency (fL/U_∞)
F_d	Actuator driving force, N
f	Actuation frequency, Hz
h	Orifice slot width, m
i	Actuator current, A

j	Imaginary unit ($j = \sqrt{-1}$)
K	Nozzle non-linear damping coefficient or stiffness, N/m
L	Length, m or inductance, H
L_0	Synthetic jet stroke length, m
M	Mass, kg
p_c	Cavity acoustic pressure, Pa
R	Resistance, Ohm
Re	Reynolds number ($Re = U_0 d_h / \nu$)
s	Laplace variable ($s = j2\pi f$), s^{-1}
t	Time, s
U	Velocity, m/s
U_0	Characteristic velocity ($U_0 = L_0 f$), m/s
V	Volume, m^3
x	Displacement, m
Z	Impedance, Ω

Subscripts

0	Characteristic synthetic jet scale
c	Actuator cavity
d	Actuator diaphragm
e	Electrical (referring to the actuator driver)
emf	Electromotive force
H	Helmholtz resonance
n	Actuator nozzle or orifice
p	Parallel
s	Series
∞	Free stream

Greek symbols

α	Electromechanical coupling coefficient, N/V
β	Nozzle added mass coefficient
ρ	Density, kg/m^3
ν	Kinematic viscosity, m^2/s

Abbreviations

CFD	Computational fluid dynamics
ROM	Reduced-order model
SJA	Synthetic jet actuator

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